

# Directly probing the mechanical properties of the spindle and its matrix

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Several recent models for spindle length regulation propose an elastic pole to pole spindle matrix that is sufficiently strong to bear or antagonize forces generated by microtubules and microtubule motors. We tested this hypothesis using microneedles to skewer metaphase spindles in *Xenopus laevis* egg extracts. Microneedle tips inserted into a spindle just outside the metaphase plate resulted in spindle movement along the interpolar axis at a velocity slightly slower than microtubule poleward flux, bringing the nearest pole toward the needle. Spindle velocity decreased near the pole, which often split apart

slowly, eventually letting the spindle move completely off the needle. When two needles were inserted on either side of the metaphase plate and rapidly moved apart, there was minimal spindle deformation until they reached the poles. In contrast, needle separation in the equatorial direction rapidly increased spindle width as constant length spindle fibers pulled the poles together. These observations indicate that an isotropic spindle matrix does not make a significant mechanical contribution to metaphase spindle length determination.

## Introduction

Proper assembly of meiotic/mitotic spindles at metaphase yields bipolar structures with cell type-specific steady-state lengths. This spindle morphology is a prerequisite for accurate chromosome segregation. Despite the relatively rapid turnover of its dynamic microtubule building blocks, the spindle achieves a steady-state length that is determined by the concerted action of microtubule dynamics, motors, and other spindle-associated proteins (Walczak and Heald, 2008). How this self-assembly process is achieved is still not fully understood, but it is clear that forces play a critical role.

The current thinking is that steady-state spindle length is determined by a force-balance mechanism in which outward-directed forces that push the spindle poles apart are antagonized by inward-directed forces that pull them together (Mogilner et al., 2006). In higher eukaryotes, the plus end-directed kinesin-5 (Eg5 in *Xenopus laevis*) is thought to slide oppositely oriented microtubules that overlap in the spindle midzone in a microtubule minus end direction toward opposite poles. This

produces persistent poleward movement of microtubules within each spindle half, which is called flux (Sawin and Mitchison, 1991; Miyamoto et al., 2004), and generates an outward pushing force that works to keep the spindle poles spatially separated (Sharp et al., 1999; Mitchison et al., 2005). Minus end-directed motor complexes such as dynein–dynactin and kinesin-14 family members have been shown to produce forces on spindle poles that antagonize kinesin-5 in determining metaphase spindle length (Mitchison et al., 2005; Tanenbaum et al., 2008; Ferenz et al., 2009). It is clear that these motor-driven pushing and pulling forces depend on microtubules, but an elastic, nonmicrotubule spindle matrix has also been proposed to be an important mechanical component of the spindle (Pickett-Heaps and Forer, 2009; for review see Scholey et al., 2001). Indeed, an elastic matrix of sufficient strength could mechanically influence spindle length by opposing forces generated within the spindle by microtubule-dependent motors (Goshima et al., 2005; Mitchison et al., 2005) or by serving as a scaffold

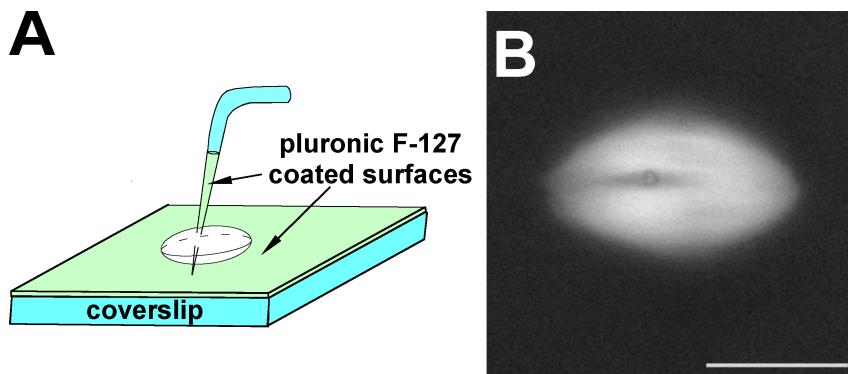
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Abbreviations used in this paper: CSF, cytostatic factor; DIC, dynein intermediate chain; NuMA, nuclear mitotic apparatus protein.

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**Figure 1. Experimental approach for spindle skewering.** (A) The cartoon shows the experimental setup used in all skewering experiments (see Materials and methods for more details). (B) Example of a skewered spindle visualized by the addition of X-rhodamine-labeled tubulin to the extract. The cross section of the microneedle is seen as a dark annulus. Bar, 25  $\mu$ m.



for kinesins and dyneins that slide microtubules (Kapoor and Mitchison, 2001; Tsai et al., 2006).

Although several spindle matrix proteins/molecules have been identified (Walker et al., 2000; Chang et al., 2004; Qi et al., 2004; Tsai et al., 2006; Fabian et al., 2007), there is still uncertainty in the spindle mechanics field about whether these spindle matrices make mechanical contributions to the determination of spindle length. In this study, we use microneedle manipulation to directly probe the mechanics of metaphase spindles in *Xenopus* egg extracts. Our findings not only address issues about the mechanical properties and structural organization of the spindle matrix but also have implications for the distribution and mechanical properties of lateral linkages between microtubules within the spindle.

## Results and discussion

### Intrinsic forces push skewered spindles off stationary microneedles

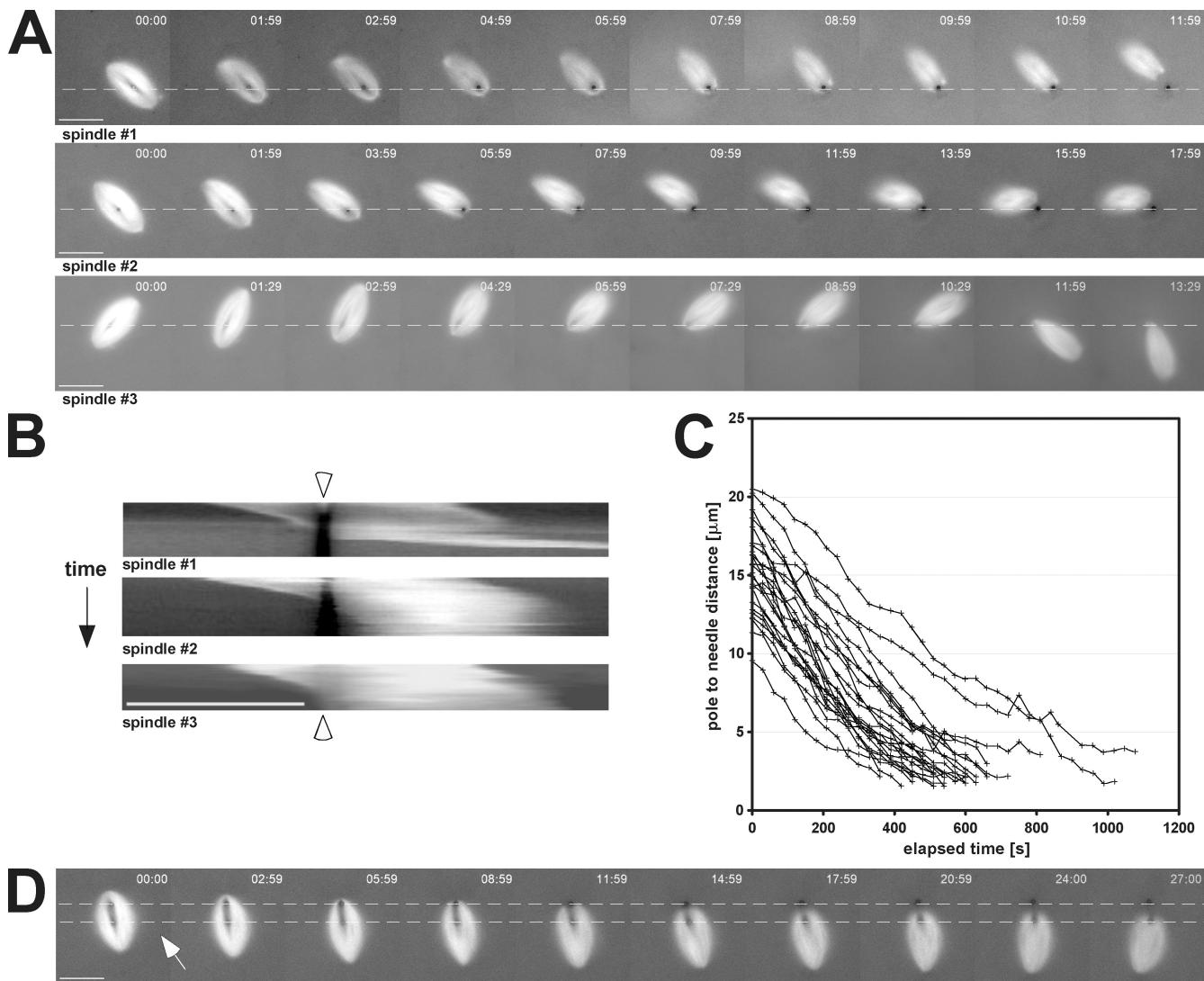
We assembled meiosis II spindles in *Xenopus* egg extracts as described previously (Murray, 1991; Gatlin et al., 2009) and visualized them by the addition of trace X-rhodamine-labeled tubulin or polarization microscopy. A small aliquot of extract with assembled spindles was spread on a coverslip surface passivated with a hydrophobic film to minimize adventitious protein binding, covered with a thin layer of mineral oil to prevent evaporation, and mounted on an inverted microscope. This open setup allowed access for two micromanipulator-controlled microneedles. The needles were bent near their tips at an  $\sim 70^\circ$  angle so that their tips could be pushed vertically through a spindle, pinning it to the coverslip (Fig. 1 A). The needle tips were also passivated.

In initial experiments, a single needle was pushed through a spindle at a point lying on the spindle interpolar axis and just poleward of its equator to avoid entanglements with chromosome arms (Fig. 1 B). Near the initial needle insertion point, more microtubules are oriented with their plus ends toward the equator than the other way around, and thus the majority of microtubules slide toward the pole closest to the needle (Burbank et al., 2006; Yang et al., 2008). As the spindle fiber microtubules flux poleward, molecular cross-links between adjacent microtubules, if present and unbroken, would be expected to move with the fluxing microtubules. Such cross-links would push against the fixed needle and move

the whole spindle in the opposite direction if not constrained by a spindle matrix.

In 24 out of 27 cases, the spindle pushed itself along its interpolar axis, in the expected direction, until it moved all of the way off the microneedle (Fig. 2 A and [Video 1](#)). In the three cases in which the spindle did not move past the needle, the spindle was skewered very close to the exact middle of the spindle and likely became entangled in chromatin. Inhibition of microtubule sliding using AMP-PNP (Sawin and Mitchison, 1991) blocked movement of the spindle relative to the skewering needle (Fig. 3 C), suggesting that movement of the spindle relative to the needle is likely caused by the same forces that produce microtubule sliding, presumably kinesin-5 activity (Miyamoto et al., 2004). Spindles sometimes rotated around the microneedles (e.g., Fig. 2 A, spindles 2 and 3) because of transient and changing flows within the extracts on the coverslip. This rotation allowed the interpolar axis of the spindle to align in the direction of the extract flow, thereby increasing the contribution of viscous drag forces to spindle movement along the same axis. We attempted to reduce the influence of these forces on spindle movement by skewering spindles with two microneedles, each inserted through the same spindle half along the interpolar axis. This prevented spindle rotation and relegated the contribution of the flow forces to the vector component projected on the needle to needle axis (the white arrow in Fig. 2 D indicates the most persistent direction of flow). In each case, the spindle moved along the axis formed by the two needles and then eventually pushed itself completely off both needles ( $n = 15$ ). It should also be noted that in some single needle experiments, extract flow was noticeably less pronounced, producing instances in which two separate spindles in the same field, each skewered by a single needle, moved in different directions but at similar rates ([Video 2](#)). Collectively, these data strongly suggest that the observed spindle movements are mainly the result of force-generating mechanisms within the spindle.

The rate of spindle escape off single skewering needles was measured by kymography and manual tracking. Kymographs demonstrate that the spindles did not undergo a change in length and instead moved relative to the skewering needle as a rigid body (Fig. 2 B). By manual tracking, the average mean velocity of approach of the pole to the needle was  $1.7 \pm 0.4 \mu\text{m}/\text{min}$  for all distances  $>5 \mu\text{m}$  from the pole. This means that in most of the central half-spindle, spindle translocation relative to the needle occurred with a mean velocity close to the



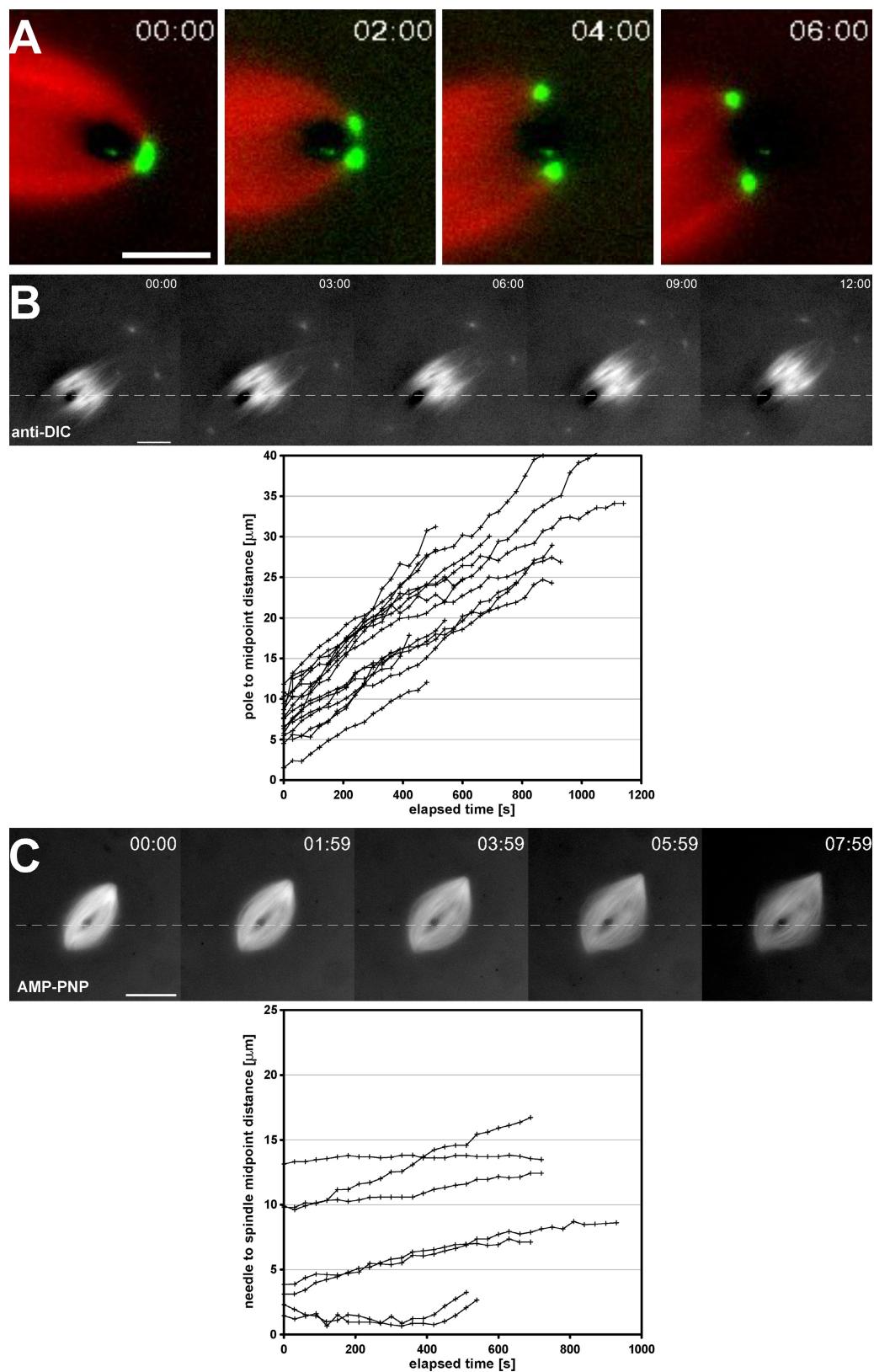
**Figure 2. Intrinsic spindle forces move impaling microneedles through the spindle.** (A) Images shown were selected from three time-lapse series of skewed spindles. Dashed lines indicate the position of the microneedle, which remained essentially stationary during the experimental time course. (B) Kymographs were used to analyze velocities of spindles as they translocated off microneedles. Because skewed spindles often rotated during time-lapse experiments as the result of changing flows within the extract (e.g., spindle #3), each image in the series was rotated using custom software to maintain a fixed orientation of the spindle's interpolar axis (arrowheads mark the position of the needle). (C) The distance between the nearest pole and the middle of the needle was plotted versus time. Slopes were calculated using linear regressions from these plots of pole to needle distances of  $>5\text{ }\mu\text{m}$  and those equal to  $5\text{ }\mu\text{m}$ . (D) Two microneedles were used to impale spindles. In each case, the needles were positioned on the same side of the spindle midzone, and the spindle moved off both needles, regardless of the predominant direction of extract flow (indicated by the white arrow). Dashed lines indicate the position of the microneedles used to skewer the spindles. Bars,  $25\text{ }\mu\text{m}$ .

range of values reported for flux velocities in *Xenopus* extract spindles ( $\sim 2\text{ }\mu\text{m}/\text{min}$  in the central region; Desai et al., 1998; Maddox et al., 2003b; Yang et al., 2007). High-resolution and confocal fluorescence speckle microscopy and computer-aided tracking were used to determine flux rates in skewed spindles (see Materials and methods). For the example in Fig. S1, the mean flux rate was similar to the measured rate of spindle translocation in the opposite direction. Although we concede that the flux-like rates of spindle movement observed do not prove a causal relationship, because we can exclude extrinsic mechanisms, it is difficult to conceive of some other intrinsic force-generating mechanism capable of producing the observed spindle movements. Thus, these observations suggest that intrinsically generated spindle forces, likely the

same responsible for producing flux, can push against a skewing microneedle.

#### Spindle translocation velocity slows as the needle encounters a pole

Within  $5\text{ }\mu\text{m}$  of the spindle pole, the velocity of spindle movement past the needle decreased to  $0.8 \pm 0.4\text{ }\mu\text{m}/\text{min}$ , which is significantly slower than for most of the half-spindle ( $P < 0.05$  by Student's *t* test for equal means;  $n = 24$ ). The rate of microtubule sliding near the pole is currently controversial; in a single molecule study, sliding rates at  $\sim 80\%$  of the equatorial value were reported (Yang et al., 2008), whereas a study using an averaging method reported that sliding slows to near zero velocity at the pole (Burbank et al., 2007). Therefore, from velocity



**Figure 3. Intrinsic forces push but do not pull the spindle off impaling microneedles.** (A) The time-lapse series shows the dynamic morphology of a spindle pole, labeled with Alexa Fluor 488 anti-NuMA antibodies (green), as it is split by a microneedle. Spindles continued to move despite a lack of any detectable microtubules (red) on the distal side of the microneedle. (B) Time-lapse images show the behavior of skewered spindles  $\sim$ 5–10 min after the addition of function-perturbing antibodies against the 70.1-kD DIC (anti-DIC). The distance between the metaphase plate and the needle were measured and plotted as a function of time for multiple spindles in the corresponding graph. (C) Assembled spindles were treated with 1.5  $\mu$ M AMP-PNP and skewered within 5–10 min after treatment. AMP-PNP at this concentration inhibited flux, in agreement with nearly horizontal plots of needle to spindle midpoint distance versus time (graph). (B and C) Dashed lines indicate the position of the microneedle used to skewer the spindles. Bars: (A) 5  $\mu$ m; (B and C) 25  $\mu$ m.

measurements alone, it is difficult to determine whether spindle translocation relative to the needle slowed in our experiments because the polar region contains some additional structural element that resists entry of the needle (i.e., a dynein–dynactin–nuclear mitotic apparatus protein [NuMA] pole matrix; Dionne et al., 1999) or because microtubule sliding is intrinsically slower at the poles. Nevertheless, almost all spindles eventually translocated off the needle by splitting the pole apart as the needle passed through it (Fig. 3 A and [Video 3](#)), suggesting that if there is a pole matrix, it must be sufficiently weak or sufficiently dynamic to eventually allow the needle to pass through it.

Dynein–dynactin complexes play a major role in organizing the poles of egg extract spindles, probably by acting as a cross-linker themselves, but also by transporting other cross-linkers such as NuMA to the pole (Heald et al., 1996; Merdes et al., 2000; Gaetz and Kapoor, 2004; Mitchison et al., 2005). To test the role of dynein–dynactin in slowing needle passage through the pole, we perturbed its function by the addition of an inhibitory antibody to the 70.1-kD dynein intermediate chain (DIC; ab70.1 at 1 mg ml<sup>-1</sup>; Heald et al., 1996). As expected, this treatment caused poles to splay out and spindle length to increase (Fig. 3 B). Like untreated control spindles, skewered dynein-inhibited spindles with unfocused poles moved off the needle (Fig. 3 B and [Video 4](#)). The escape velocity was determined by measuring the change over time in the distance between the needle and spindle midpoint. This speed was similar to the value measured in the central half-spindle region for untreated spindles ( $1.6 \pm 0.4$ ;  $n = 17$ ;  $P = 0.4$ ). The velocity was relatively constant even as separation distances approached the typical length of untreated half-spindles,  $\sim 17$   $\mu\text{m}$ . Thus, dynein–dynactin function is necessary for producing the much slower spindle velocities relative to the needle within the pole region of control spindles.

#### Spindle deformation using two needles

To qualitatively assess regional differences in spindle mechanical properties, we used two skewering needles, moved them apart along either spindle axis, and observed the effects on spindle morphology.

**Longitudinal (pole to pole) movement.** When needles were inserted one on either side of the spindle equator and then moved along the interpolar axis at rates faster than flux velocity, spindles were initially unchanged in length or morphology. Only when both needles reached the poles did the spindle stretch (Fig. 4, A and B; and [Video 5](#);  $n = 11$  spindles). Increases in spindle length only began when the needles were separated by distances greater than  $\sim 80\%$  of the original spindle length and stopped shortly after one of needles broke through a pole (Fig. 4 B;  $n = 4$  spindles, which remained in focus throughout pull). The relatively short time scale over which the elongation occurred and a concurrent decrease in spindle width strongly suggest that when the needles reached the pole, the spindles became deformed (i.e., stretched) as opposed to being elongated as the result of some mechanochemical change in microtubule dynamics or microtubule motor function (Dumont and Mitchison, 2009). These results reveal large differences in the mechanical properties of the spindle matrix or other microtubule

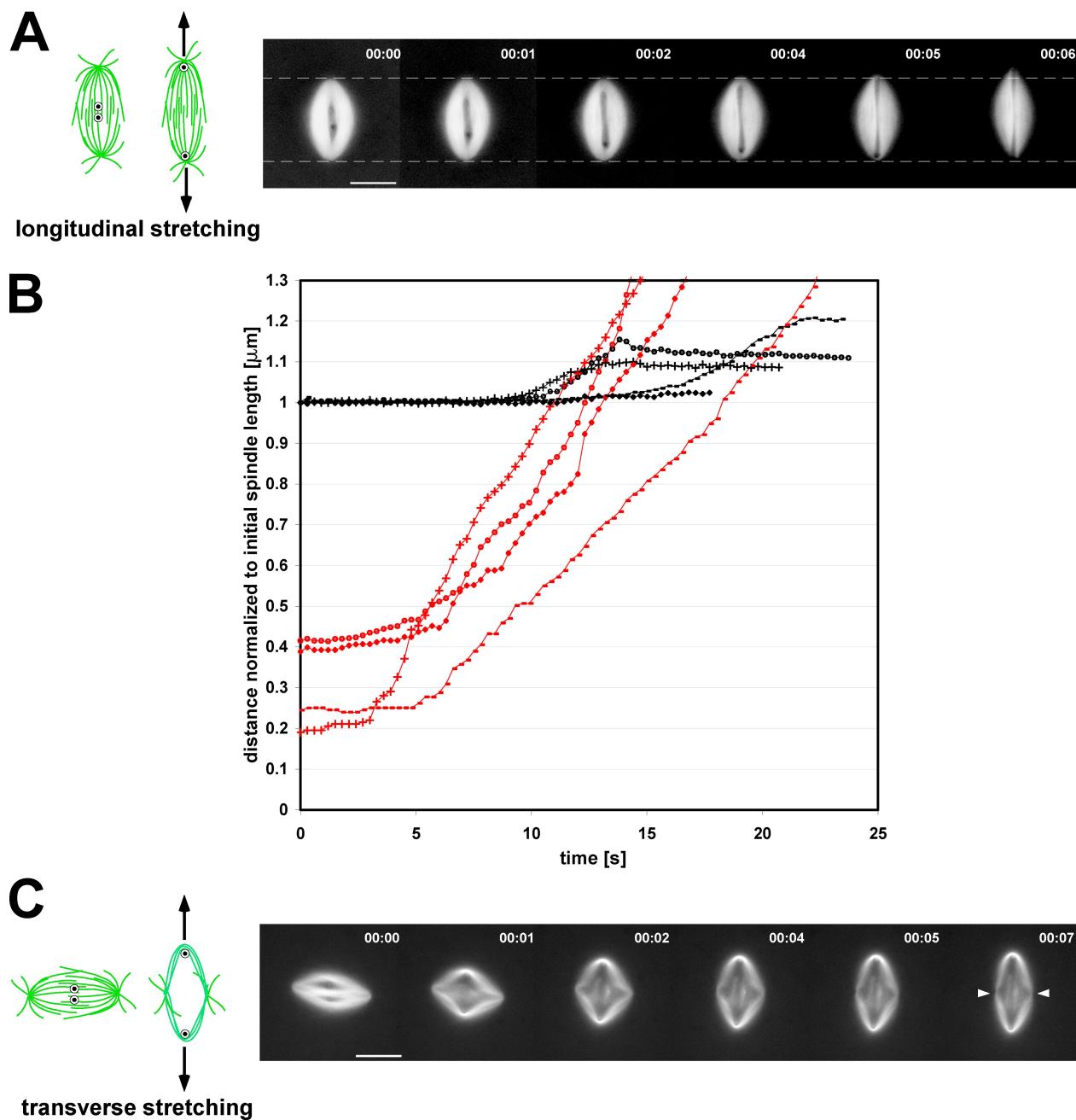
cross-linkers between the central half-spindle and the poles. Furthermore, they suggest that the slow movement of the needle through the poles in the skewering experiments is not caused by locally slower flux velocities but instead results from enhanced mechanical resistance.

**Transverse movement.** When the two needles were initially inserted along the spindle equator and subsequently moved apart perpendicular to the spindle interpolar axis, the spindle stretched in the equatorial axis. As the spindle fibers were pulled outward by the needles, the spindle poles moved inward toward each other (Fig. 4 C and [Video 6](#);  $n = 11$ ). Microtubules extending in an interpolar direction inside the two needles buckled near the spindle center. The microtubules collected on the outside of each needle remained attached almost exclusively at the spindle poles, and the fiber length between the needle and the pole was similar to the normal half-spindle length. In these experiments, needle movement was stopped before the spindle was pulled apart. Deformed spindles typically exhibited two main fates: they either would rotate and expel a needle through one of their poles or simply fall off one of the needles after the disappearance of microtubules from the outside of the needle. This fiber length was stable for a period much longer than the mean half-life of most spindle fiber microtubules, which is 100 s or less ([Video 7](#); Sawin and Mitchison, 1991).

Collectively, these two needle stretching experiments provide clear evidence that an isotropic spindle matrix and lateral linkages between microtubules within the central spindle region are weak, whereas near the poles, such linkages or a pole-associated matrix is much stronger. A similar conclusion was reached  $>40$  yr ago for kinetochore fiber microtubules in grasshopper spermatocytes mechanically probed by micro-needle manipulation of their chromosome arms. A metaphase chromosome and its attached k-fiber were easily moved laterally within the spindle without much influence on adjacent chromosomes or fibers, whereas the manipulated k-fiber remained strongly anchored at its kinetochore and pole (Nicklas and Staehly, 1967; Begg and Ellis, 1979; Nicklas et al., 1982).

#### How microtubule poleward flux might push the spindle off a skewering needle

There is a long list of motor and nonmotor proteins that could potentially cross-link microtubules throughout the spindle (Walczak and Heald, 2008). We hypothesize that these cross-links, although weak, allow flux forces to work against the needle in our experiments (Fig. S2). At a spindle velocity of 1.7  $\mu\text{m}/\text{min}$ , the viscous resistance from the extract is  $<1$  pN (see Materials and methods). Thus, even weak lateral linkages between microtubules would be able to facilitate spindle movement at the velocities observed, provided other structural elements within the spindle produce little additional resistance to needle movement through the spindle. Because the lateral linkages are weak in the central region of the spindle, it is not surprising that the spindle velocity, on average, is slightly slower than the flux velocity as the result of linkage breakage and some slippage of microtubule flux past the needle. We propose that near the pole, the situation becomes different because resistance to needle movement becomes much higher.

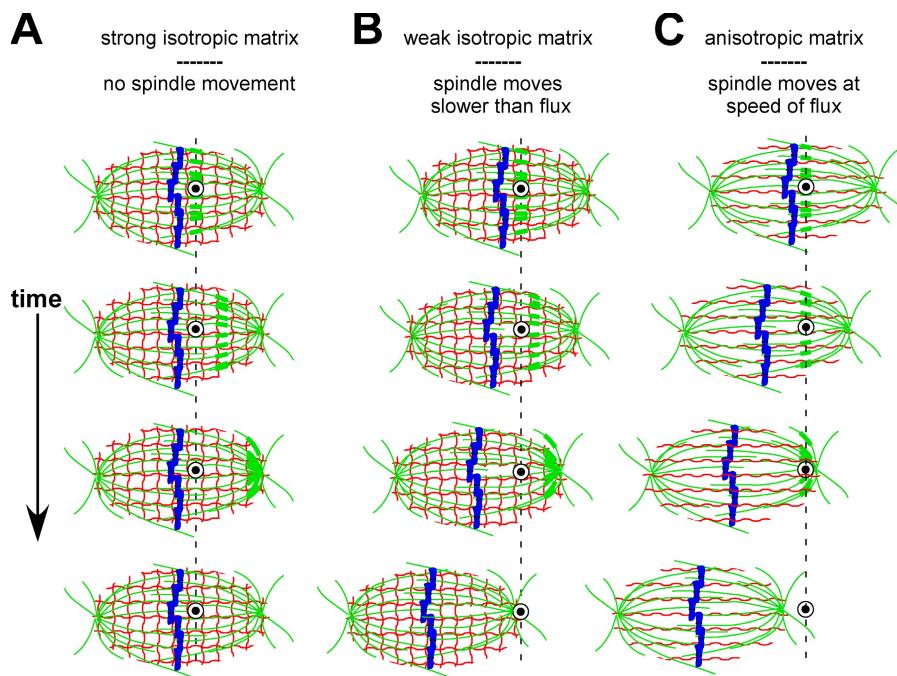


**Figure 4. Lateral interactions between microtubules are more robust near the spindle poles.** (A and C) The cartoons show the initial positions of the two skewering needle tips within the spindle and the direction of needle movement. (A) The two needles were initially positioned one on each side of the metaphase plate and then spread apart along the interpolar axis. The images show a representative time-lapse series of a spindle being longitudinally stretched. Dashed lines represent the position of the spindle poles before the onset of stretching. (B) Changes in spindle length during longitudinal stretching (normalized to the initial length of the spindle) are plotted versus time, shown as black lines, whereas corresponding plots of needle separation versus time are shown in red. Matching markers indicate data taken from the same experiment. Percentages are the ratio of needle separation at the onset of spindle elongation to the initial spindle length (elongation onset was arbitrarily defined as the time point at which the spindle became 0.5% longer than its initial length). (C) Spindles were also stretched in the orthogonal direction, transverse to their interpolar axes. In these experiments, spindle deformation began at the onset of needle separation and continued until the needles were stopped. White arrowheads indicate the position of the interpolar axis after transverse spindle stretching. Bars, 25  $\mu\text{m}$ .

#### Does a spindle matrix play a mechanical role in spindle length regulation?

The rationale for our single needle skewering experiments was based on predictions about resultant spindle movement and resistance to this movement by spindle matrices of various structure and material properties (Fig. 5). If the spindles did not move

at all, it would have been difficult to distinguish between (a) the absence of a force-transmitting linkage between fluxing microtubules and the needle or (b) the presence of an isotropic matrix of sufficient mechanical strength to completely uncouple needle movement from flux-driven forces (Fig. 5 A). In contrast, the spindles did move through, and eventually off, the



**Figure 5. Predicted experimental outcomes of single needle experiments for different types of matrices assuming flux forces can be transmitted to the needle by lateral microtubule cross-links.** (A–C) Thin red lines represent the matrix, whereas green lines represent spindle microtubules. Thicker sections along the microtubules depict the movement of a photoactivated mark on the microtubule lattice to illustrate poleward flux over time. The blue shapes within the spindle midzone represent aligned metaphase chromosomes. The dashed line runs through the midpoint of the stationary needle tip, as seen in cross-section (depicted as a white annulus with a filled, black center). See “Does a spindle matrix play a mechanical role in spindle length regulation?” for explanations of each cartoon.

skewering needles. It is this observation that has major implications regarding the properties of the spindle matrix and distinguishes our approach from previous micromanipulation studies of the spindle (Begg and Ellis, 1979; Nicklas et al., 1982; Itabashi et al., 2009). Our data support the conclusion that if an isotropic matrix is present throughout the spindle or encases it as a sheath, it is sufficiently weak enough to allow fluxing microtubule arrays to push a skewering needle through it with a velocity at or slightly slower than the flux velocity within the central half-spindle (Fig. 5 B). Therefore, when one considers the time scale over which spindle length is maintained, it is difficult to imagine how such an isotropic matrix could make a mechanically significant contribution to the steady-state length of the metaphase spindle. However, our experiments do not exclude the possibility of a structurally anisotropic, filamentous matrix that aligns itself with spindle microtubules (Fig. 5 C). Such a matrix could be strong enough along the interpolar axis to contribute mechanically to the force balance thought to regulate spindle length yet have very weak lateral associations, if any at all, so that fluxing microtubule arrays push the needle by them within the central half-spindle at flux-like speeds.

Our data suggest that to play a relevant mechanical role in regulating spindle length, a spindle matrix must be structurally anisotropic. Although it is clear that perturbation of matrix proteins/molecules leads to aberrant spindle morphology and function (Chang et al., 2004; Tsai et al., 2006; Lince-Faria et al., 2009), it is not clear whether any of the currently identified matrix proteins/molecules are arranged in a manner consistent with this requirement (Walker et al., 2000; Chang et al., 2004; Qi et al., 2004; Tsai et al., 2006; Lince-Faria et al., 2009). Thus, in the absence of higher-resolution structural data, we can only speculate about whether these matrix components exhibit structural anisotropy, thus leaving open the possibility that each could regulate spindle length mechanically. Alternatively,

matrix proteins/molecules might play other, nonmechanical roles in the process; perhaps by acting as scaffolds for regulatory proteins that govern microtubule-dependent spindle assembly and function (Tsai et al., 2006) or as local modulators of diffusion (as suggested in Lince-Faria et al. [2009]).

Our results do support the existence of enhanced microtubule cross-linking and/or the presence of a strong isotropic matrix near spindle poles, which is eliminated by inhibiting dynein function. Microtubule cross-linking could be produced directly by dynein–dynactin–NuMA complexes, which are also known to aggregate and form an insoluble pole matrix with distinct dynamics from microtubules (Dionne et al., 1999). In addition, motor-dependent transport and concentration of other microtubule cross-linkers or potential matrix components might also contribute to the increased resistance to needle movement through the poles (e.g., Eg5; Uteng et al., 2008). These spindle components concentrated at poles might provide both the firm anchorage of kinetochore fiber microtubule minus ends, as initially shown by Nicklas and Staehly (1967), and the viscoelastic coupling between adjacent bundles of kinetochore microtubules that appears to be required for synchronous segregation of chromosomes in anaphase (Matos et al., 2009).

## Materials and methods

All reagents were purchased from Sigma-Aldrich unless otherwise stated.

### Spindle assembly in *Xenopus* egg extracts

Cytostatic factor (CSF)-arrested egg extracts were prepared as described previously (Murray, 1991; Desai et al., 1999). In brief, to assemble meiosis II spindles with replicated chromosomes, extracts were induced to go into interphase by the addition of excess  $\text{Ca}^{2+}$  and demembranated sperm nuclei (200 nuclei/ $\mu\text{l}$  final concentration). As an additional step to improve the consistency of extracts for micromanipulation experiments, we used a syringe filter (1.2- $\mu\text{m}$  pore size) to remove residual pigment granules and any large, insoluble aggregates. Compared with spindles assembled in unfiltered extracts, this treatment had no discernible effect on spindle

assembly or morphology. Extracts were then incubated at 18°C for 1 h to allow for DNA replication and driven into metaphase by the addition of an equal volume of CSF extract. To more readily visualize needles within the spindle, X-rhodamine-labeled tubulin was prepared as previously described (Yang et al., 2007) and added to extracts at a final concentration of  $\sim 5 \mu\text{g ml}^{-1}$  immediately after add-back of CSF-arrested extract to the spindle assembly reactions. Spindles were then allowed to assemble for 1–1.5 h at 18°C before experimentation.

#### Preparation of coverslips

Our skewering experiments were performed on or near the coverslip surface. So, to limit adventitious binding of motors and other proteins that could potentially influence the movements of skewered spindles, we acid washed our coverslips (#1.5 thickness; Corning) as described previously (Waterman-Storer et al., 1999). The coverslips were then dried under a stream of nitrogen and silanized by immersion in a solution of 0.05% dichlorodimethylsilane (Fluka) in trichloroethylene for 2 h. Silanized coverslips were rinsed in methanol with bath sonication for 5 min (three times) and then in double-distilled water for 5 min (three times) and allowed to air dry until use. Before experimentation, treated coverslips were covered in *Xenopus* buffer (Desai et al., 1999) with 1% Pluronic F-127, a hydrophobic coblock polymer (Varga et al., 2006). After at least 10 min, the *Xenopus* buffer + Pluronic buffer was removed, and the coverslip was rinsed under a stream of double-distilled water. Any residual rinse was removed by tilting the coverslip and wiping its edge using a Kim wipe.

#### Spindle micromanipulation and imaging

The microneedles used for spindle manipulation were made by pulling borosilicate glass micropipettes using a micropipette puller (model P-97; Sutter Instrument Co.). The distal tips were bent at angles ranging from 60 to 90° by heating the needle on its bottom side with a microforge filament (MF-830; Narishige). The bent needles were then silanized and coated with Pluronic F-127 as described in the previous section for coverslips.

Chambers used for all skewering experiments were prepared as described previously (Gatlin et al., 2009). In brief, a 22-mm diameter coverslip was adhered to a preheated metal slide using VALAP (1:1:1 vasoline, lanolin, and paraffin) applied around a circular cutout in the slide. This formed a shallow well with the coverslip acting as the bottom. A small aliquot ( $\sim 5\text{--}7 \mu\text{l}$ ) of extract was spread over the surface of the coverslip with a pipette tip and overlaid with 300  $\mu\text{l}$  of mineral oil (embryo tested from Sigma-Aldrich) to prevent the sample from drying out during subsequent imaging.

Micromanipulation/skewering experiments were conducted on an inverted microscope (TE2000e; Nikon) stand equipped with a motorized microscope stage (Applied Scientific Instrumentation). Dual manual micromanipulators (Narishige) were mounted directly to the condenser arm of the microscope and adjusted so the shafts of opposing needles entered the well at shallow angles with the bent needle tips nearly perpendicular to the coverslip surface. Needles were lowered through the mineral oil overlay and into the extract for skewering. Polarization optics (i.e., a crossed polarizer and analyzer) in combination with a 20x differential interference contrast objective (20x NA 0.75 M1mm; Plan Fluor; Nikon) were used to visualize spindles during skewering. For subsequent imaging of spindle escape, we used epifluorescence microscopy to view spindle-incorporated X-rhodamine-tubulin (as prepared in Waterman-Storer et al. [1999]). Where indicated, pole localization of NuMA was visualized using anti-NuMA antibodies directly labeled with Alexa Fluor 488 (Invitrogen; a gift from A. Groen, Yale University, New Haven, CT). Images were acquired using a cooled charge-coupled device camera (Orca-ER; Hamamatsu Photonics) mounted on the right-hand side of the microscope stand to a port that received 80% of the emitted light, with the remaining 20% going to the eyepieces. This setup allowed us to manipulate and image spindles during transverse and longitudinal stretching experiments. These experiments were conducted at  $\sim 20^\circ\text{C}$ . All image acquisition and analyses were performed using MetaMorph software (MDS Analytical Technologies). For kymograph analysis of spindle escape, rotation of acquired image stacks was performed using custom Matlab software written by J.C. Gatlin. Unless otherwise stated, at least three different extracts were used for each experimental condition.

#### Statistical analyses

Excel software (Microsoft) was used for all statistical analysis and graphing. Unless otherwise indicated, reported results are expressed as mean values  $\pm$  standard deviations.

#### Calculations of viscous drag force on spindles in *Xenopus* egg extracts

The viscous force resisting spindle movements in the *Xenopus* egg extract is dependent on the viscosity of the extract, the geometry of the microtubule array, and the velocity at which the array is moved. At low Reynolds numbers, the drag forces acting on a spindle modeled as a solid, prolate ellipsoid and moving at a velocity  $u$  in a fluid of viscosity  $\eta$  is

$$F_{\parallel} = 4 \frac{\pi \eta r u_{\parallel}}{[\ln(2r/h) - 1/2]},$$

where  $r$  is the longest semiaxis of the ellipse and  $h$  is the shortest semiaxis of the ellipse (Perrin, 1934; Berg, 1983).

The viscosity of crude interphase extracts has been measured at  $\sim 20 \text{ mPa}\cdot\text{s}$  (Valentine et al., 2005), and the velocities of spindle movements in these experiments are in the order of  $\sim 2 \mu\text{m}/\text{min}$ . For a typical spindle in our experiments,  $r/h$  is  $\sim 2$  with an  $r$  of  $\sim 17 \mu\text{m}$ . Substituting these numbers into the aforementioned equation for  $F_{\parallel}$  (i.e., spindles sliding parallel to their long axes) yields  $F_{\parallel} \equiv 0.2 \text{ pN}$ .

#### Determining the rate of flux in a skewered spindle

We measured the rate of spindle translocation relative to a stationary skewering needle by manually tracking the nearest pole and calculating the distance between it and the needle as a function of time (Fig. S1 A). This velocity,  $V_{\text{spindletrans}}$ , was  $2.3 \mu\text{m}/\text{min}$  determined by linear regression ( $R^2 = 0.99$ ). Confocal fluorescent speckle microscopy (Maddox et al., 2003a) was used to image microtubule flux within the skewered spindle on a microscope (Eclipse TE300; Nikon; differential interference contrast 60x NA 1.4). Acquisition was controlled with MetaMorph software, and all images were acquired at room temperature ( $\sim 22^\circ\text{C}$ ). A computer-based tracking method that utilizes an optimal-flow minimum-cost correspondence assignment was used to track speckle flows (Yang et al., 2005). Four main directional components (along the interpolar and equatorial axes of the spindle) and the corresponding vectors for each were extracted, insuring that velocity contributions from spindle rotations were isolated from speckle flows along the interpolar axis. The results presented in Fig. S1 B are for those flows running parallel to the pole to pole axis and directed toward the pole farthest from the needle. The mean velocity of these speckles,  $V_{\text{speckles}}$ , was measured as  $4.6 \pm 1.4 \mu\text{m}/\text{min}$  ( $n = 189$  tracks; Fig. S1 C). Because the spindle is moving in the same direction as this flux  $V_{\text{speckles}} = V_{\text{flux}} + V_{\text{spindletrans}}$ . Substituting our measured values of  $V_{\text{speckles}}$  and  $V_{\text{spindletrans}}$ ,  $V_{\text{flux}}$  is  $2.3 \pm 1.4 \mu\text{m}/\text{min}$ . In this case, the spindle translocates off the needle at the same mean speed of its flux.

#### Online supplemental material

Fig. S1 shows the computer-aided tracking of microtubule flux in a skewered metaphase spindle. Fig. S2 shows a model for how flux pushes against the needle. Video 1 shows time-lapse fluorescence imaging of spindle escape off a single skewering needle. Video 2 shows time-lapse fluorescence imaging of two spindles in the same frame escaping their needles in different directions. Video 3 shows time-lapse fluorescence imaging of a skewering needle moving through a spindle pole. Video 4 shows time-lapse fluorescence imaging of spindle escape after treatment with concentrated anti-DIC antibodies. Videos 5 and 6 show time-lapse fluorescence imaging of two-needle separation along the interpolar spindle axis (longitudinal stretching) and orthogonal to the interpolar spindle axis (transverse stretching), respectively. Video 7 shows time-lapse fluorescence imaging of spindle behavior after a transverse stretch. Online supplemental material is available at <http://www.jcb.org/cgi/content/full/jcb.200907110/DC1>.

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